

## ON THE STIFFNESS PROPERTY OF MOTION

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It is well known that the axis of a rapidly rotating gyroscope is very little responsive to large perturbing forces, i. e., possesses stiffness [1]. It was noticed that under specific conditions the stiffness property is inherent in systems with gyroscopes [2]. Many mechanical systems possess a property skin in some sense to gyroscopic stiffness. In this paper stiffness is interpreted as a distinctive stability. Theorems establishing tests for stiffness, analogous to the theorems of Liapunov's direct method, are formulated. The stiffness property is described by using a separation of variables as is done in problems on stability with respect to a part of the variables (see [3, 4], etc.). Individual questions on the stiffness of motion were examined in [5, 6].

**1. Basic definitions.** The equations of motion of a mechanical system are written as

$$dy/dt = Y(t, y, g) \quad (1.1)$$

where  $t \geq 0$  is time,  $y$  is an  $n$ -dimensional state vector of the system and  $g$  is a constant vector-valued physical parameter. The motions (partial solution of (1.1))

$$y = f(t, y_0, g) \quad (1.2)$$

from some family, satisfying the initial conditions:  $y = y_0$  when  $t = t_0$ , are considered to be the unperturbed motions. It should be noted that the values of  $y_0$  and of parameter  $g$  can be related by the existence conditions for the motions (1.2).

Setting  $y = f + x$  in (1.1), we obtain the equation of perturbed motion

$$dx/dt = Y(t, f + x, g) - Y(t, f, g) \quad (1.3)$$

in which  $y_0$  and  $g$  are parameters. Equation (1.3) is considered dependent upon parameters  $a_1, \dots, a_r$  essential to the motion stiffness problem and is written as

$$\begin{aligned} dx/dt &= X(t, x, a), \quad X(t, 0, a) \equiv 0 \\ a &= (a_1, \dots, a_r) \end{aligned} \quad (1.4)$$

We shall investigate the stiffness of the motion  $x = 0$  with respect to a part of the variables  $x_\alpha$  ( $\alpha = 1, \dots, m$ ;  $m < n$ ). We assume that the vector-valued function  $X(t, x, a)$  in (1.4) is continuous and satisfies the uniqueness conditions for the solution in the domain

$$\Gamma_0 = \{x: |x_\alpha| < H, |x_\beta| < \infty \ (\alpha = 1, \dots, m; \beta = m + 1, \dots, n)\}, \quad t \geq t^* \quad (1.5)$$

for all  $a \in D$ , where  $D$  is some domain in the space of the parameters being examined. We denote the solution of (1.4) by  $x(t, t_0, x_0, a)$  and in the space of variables  $\{x_1, \dots, x_n\}$  we introduce a parallelepiped defined by the inequalities

$$\bar{\Pi}_\varepsilon = \{x : |x_\alpha| \leq \varepsilon_1, |x_\beta| \leq \varepsilon_2\}, \quad \bar{\Pi}_\delta = \{x : |x_\alpha| \leq \delta_1 < \varepsilon_1, |x_\beta| \leq \delta_2 < \varepsilon_2\} \tag{1.6}$$

the set of boundary points is denoted  $\bar{\Pi} \setminus \Pi$  in what follows.

**Definition 1.** Motion  $x = 0$  possesses stiffness with respect to variables  $x_\alpha$  if for any numbers  $\varepsilon_1 > 0$  and  $\delta_2 > 0$  (the first can be arbitrarily small, the second, arbitrarily large) and for an instant  $t_0 \geq t^*$  we can find a parameter  $a^* \in D$  and numbers  $\varepsilon_2 > 0$  and  $\delta_1 > 0$  defining domain (1.6), for which  $x(t, t_0, x_0, a^*) \in \Pi$  when  $t \geq t_0$  if only  $x_0 \in \bar{\Pi}_\delta$ . The stiffness is said to be uniform in  $t_0$  if  $a^*, \varepsilon_2$  and  $\delta_1$  do not depend on  $t_0$ .

**Definition 2.** Motion  $x = 0$  possesses strong stiffness with respect to variables  $x_\alpha$ , and domain  $\bar{\Pi}_\delta$  lies in its domain of attraction, if it possesses stiffness with respect to these variables and, in addition, the condition

$$\lim_{t \rightarrow \infty} x(t, t_0, x_0, a^*) = 0, \quad x_0 \in \bar{\Pi}_\delta$$

holds.

If in Definition 1 we set  $\varepsilon_1 = \varepsilon_2 = \varepsilon$  and  $\delta_1 = \delta_2 = \delta$ , we obtain a property of the motion, close to the practical stability discussed in [7]. Using this terminology, we arrive at the following definition; we denote  $\|x\| = \max(|x_i|)$ .

**Definition 3.** Motion  $x = 0$  possesses practical stability if for any  $\varepsilon > 0$  and instant  $t_0 \geq t^*$  we can find a parameter  $a^* \in D$  and a number  $\delta > 0$  for which  $\|x(t, t_0, x_0, a^*)\| < \varepsilon$  when  $t \geq t_0$  if only  $\|x_0\| \leq \delta$ .

A motion's stiffness and its stability (also with respect to a part of the variables) are, in general, independent properties of the motion.

**Example 1. Stiffness of a gyroscope's axis in the Euler case.**

As is customary let  $Oxyz$  be the principal axes of inertia at the fastening point,  $A = B$  and  $C$  are the corresponding moments of inertia. To describe the motion we use the variables  $\gamma_1, \gamma_2, \gamma_3, p, q$  and  $r$ , where the  $\gamma_i$  are the cosines of the angles formed by some fixed axis  $z_1$  with the axes  $x, y$  and  $z$ , respectively. In the unperturbed motion we set

$$\gamma_1^\circ = \gamma_2^\circ = 0, \quad \gamma_3^\circ = 1, \quad p^\circ = q^\circ = 0, \quad r^\circ = \omega = \text{const} \tag{1.7}$$

In the perturbed motion we denote

$$\begin{aligned} \gamma_1 &= \eta_1, \quad \gamma_2 = \eta_2, \quad \gamma_3 = 1 - \eta_3 \quad (\eta_3 \geq 0) \\ p &= \xi_1, \quad q = \xi_2, \quad r = \omega + \xi_3, \quad (\xi_3 = \text{const}) \end{aligned} \tag{1.8}$$

From the relations  $\gamma_1^2 + \gamma_2^2 + \gamma_3^2 = 1$  we can deduce that

$$\eta_1^2 + \eta_2^2 + \eta_3^2 = 2\eta_3 \tag{1.9}$$

We shall examine stiffness with respect to the variables  $\eta_\alpha$  ( $\alpha = 1, 2, 3$ ), taking  $a = \omega$ . Turning to Definition 1 we note that regarding relation (1.9) it is convenient to use the domains  $(\beta-1, 2, 3)$

$$\bar{Q}_e = \{x: \eta_1^2 + \eta_3^2 \leq \varepsilon_1^2 < 1, \eta_3 \leq 1 - \sqrt{1 - \varepsilon_1^2} < \varepsilon_1; |\xi_\beta| \leq \varepsilon_2\} \quad (1.10)$$

instead of (1.6). We specify the numbers  $\varepsilon_1 < 1$  and  $\delta_2$  and we first consider the perturbed motions under the initial conditions

$$\eta_{\alpha 0} = 0, \quad |\xi_{\beta 0}| \leq \delta_2 \quad (1.11)$$

In the perturbed motion the gyroscope's axis  $z$ , which coincides with the axis  $z_1$  at the initial instant, describes a circular cone around the vector  $L$  of the moment of momentum; the angle at the cone's apex is  $2\theta$ , where

$$\sin \theta = A\omega_{10} / L, \quad \omega_{10} = (\xi_{10}^2 + \xi_{20}^2)^{1/2}, \quad L = (A^2\omega_{10}^2 + C^2r^2)^{1/2}$$

Hence  $\cos 2\theta \leq \gamma_3(t) \leq 1$ ; consequently (see (1.8)),  $\eta_3 \leq 1 - \cos 2\theta$  or

$$\eta_3 \leq 2(A\omega_{10} / L)^2 \quad (1.12)$$

We now select the magnitude  $|\omega|$  to make the condition  $\eta_1^2 + \eta_2^2 < \varepsilon_1^2$  hold for the perturbed motions from (1.11) being examined; then  $\eta_3 < 1 - \sqrt{1 - \varepsilon_1^2}$ . Bearing (1.12) in mind, we set

$$2(A\omega_{10} / L)^2 < 1 - \sqrt{1 - \varepsilon_1^2}$$

whence follows the inequality

$$|\omega| > \delta_2(1 + \sqrt{2}A\varepsilon_1 / C(1 - \sqrt{1 - \varepsilon_1^2})) \quad (1.13)$$

From the constancy of the gyroscope's kinetic energy it follows that  $|\xi_\beta| < \varepsilon_2$  if  $\varepsilon_2 > \sqrt{2}\delta_2$ . Thus the inequalities  $\eta_1^2 + \eta_2^2 < \varepsilon_1^2$  and  $|\xi_\beta| < \varepsilon_2$  hold for the perturbed motions with initial conditions (1.11) when  $t > t_0$  if the parameter's value satisfies (1.13). However, we can be persuaded that  $\delta_1$  exists, dependent on the parameter's value chosen by (1.13), so small that the inequality holds even for  $\eta_{10}^2 + \eta_{20}^2 \leq \delta_1^2$  and  $|\xi_{\beta 0}| \leq \delta_2$ .

**2. Application of the Liapunov function method to the motion stiffness problem.** 1°. We examine real single-valued functions  $v = v(t, x, b)$  which, in general, depend on a part  $b = (a_1, \dots, a_p)$ ,  $p \leq r$ , of the variables contained in (1.4). We assume that functions  $v$  have been defined and are continuous together with the derivatives  $\partial v / \partial t$  and  $\partial v / \partial x_s$  ( $s = 1, \dots, n$ ) in the domain

$$\Gamma = \{x: |x_\alpha| \leq h < H, |x_\beta| < \infty\} \subset \Gamma_0, \quad t \geq t^*, \quad b \in D_1 \quad (2.1)$$

**Definition 4.** Function  $v(x, b)$  possesses property (A) with respect to  $x_\alpha$  if for any  $\varepsilon_1 \in (0, h)$  and  $\delta_2 > 0$  we can find  $b^* \in D_1$ ,  $\varepsilon_2 > 0$  and  $\delta_1 > 0$ , depending on them for which

$$\inf [v(x, b^*): x \in \bar{\Pi}_e \setminus \Pi_e] > \sup [v(x, b^*): x \in \bar{\Pi}_\delta] \quad (2.2)$$

Function  $v(t, x, b)$  possesses property (A) if in domain (2.1)

$$v(t, x, b) \geq w(x, b) \quad (2.3)$$

and for every  $\varepsilon_1, \delta_2$  and  $t_0 \geq t^*$  there are  $b^*, \varepsilon_2$  and  $\delta_1$  for which

$$\inf [w(x, b^*) : x \in \bar{\Pi}_\varepsilon \setminus \Pi_\varepsilon] > \sup [v(t_0, x, b^*) : x \in \bar{\Pi}_\delta] \quad (2.4)$$

Function  $v(t, x, b)$  possesses property (A) uniformly in  $t \in [t^*, \infty)$  if in domain (2.1)

$$W(x, b) \geq v(t, x, b) \geq w(x, b) \quad (2.5)$$

and for every  $\varepsilon_1$  and  $\delta_2$  there are  $b^*, \varepsilon_2$  and  $\delta_1$  for which

$$\inf [w(x, b^*) : x \in \bar{\Pi}_\varepsilon \setminus \Pi_\varepsilon] > \sup [W(x, b^*) : x \in \bar{\Pi}_\delta] \quad (2.6)$$

Having denoted  $r_1 = \max(|x_\alpha|)$  and  $r_2 = \max(|x_\beta|)$ , we consider the domain  $G_\varepsilon = \{x : 0 \leq r_1 \leq \varepsilon, 0 \leq r_2 < \infty\}$  and its boundary point set  $R_\varepsilon = \{x : r_1 = \varepsilon, 0 \leq r_2 < \infty\}$  for some  $\varepsilon \in (0, h)$ .

**Lemma.** For the function  $v(x, b)$  to possess property (A) with respect to  $x_\alpha$  it is sufficient that the following conditions be fulfilled:

- a)  $v = v^*(x_\beta)$  if  $x_\alpha = 0$ ;
- b) for any arbitrarily small  $\varepsilon > 0$  and large  $M > 0$  there exists a parameter  $b^* \in D_1$  for which  $\inf v(x, b^*) > M$  on the set  $x \in R_\varepsilon$ ;
- c)  $v(x, b) \rightarrow +\infty$  as  $r_2 \rightarrow +\infty$  uniformly relative to  $x_\alpha$  in domain  $G_\varepsilon$ .

**Proof.** Let us show that domains  $\bar{\Pi}_\varepsilon$  and  $\bar{\Pi}_\delta$  for which inequality (2.2) is fulfilled can be constructed for functions satisfying the Lemma's hypotheses. Indeed, let  $\varepsilon_1$  and  $\delta_2$  be specified. Then

$$\sup [v(x, b) : x_\alpha = 0, |x_\beta| \leq \delta_2] = M(\delta_2)$$

We select  $b^* \in D_1$  as to have

$$\inf v(x, b^*) > M(\delta_2), \quad x \in R_{\varepsilon_1}$$

We can find  $\varepsilon_2 > \delta_2$  large enough to make

$$\inf [v(x, b^*) : x \in \bar{\Pi}_{\varepsilon_2} \setminus \Pi_{\varepsilon_2}] > M(\delta_2)$$

Finally, because the function  $v(x, b^*)$  is uniformly continuous, a sufficiently small  $\delta_1$  defining  $\bar{\Pi}_\delta$  in (2.2) exists in the domain  $\bar{\Pi}_{\varepsilon_2}$ .

Function  $v(t, x, b)$  obviously possesses property (A) if  $v = v^*(t, x_\beta)$  when  $x_\alpha = 0$  and a function  $w(x, b)$  satisfying (2.3) and the Lemma's hypotheses exist. If functions  $W(x, b)$  and  $w(x, b)$  satisfying (2.5) and the Lemma's hypotheses exist, then function  $v(t, x, b)$  possesses property (A) uniformly in  $t \in [t^*, \infty)$ .

**Note 1.** We set  $\varepsilon_1 = \varepsilon_2 = \varepsilon$  and  $\delta_1 = \delta_2 = \delta$ . We say that a function  $v(x, b)$  that for this case satisfies condition (2.2) in Definition 4 possesses property (B).

According to the Lemma we can conclude that  $v(x, b)$  possesses property (B) if  $v(0, b) \equiv 0$  and for every  $\varepsilon$  (no matter how small) there exists  $b^*$  for which  $\inf [v(x, b^*) : \|x\| = \varepsilon] > 0$ .

2°. Some indications of motion stiffness. Theorem 1. If for system (1.4):

- a) there exists a function  $v(t, x, b)$ , possessing property (A) with respect to  $x_\alpha$ ;
- b) function  $v$  and its derivative  $v'$  (by virtue of system (1.4)) satisfy the condition: for every  $\varepsilon_1, \delta_2$  and  $t_0 \geq t^*$  specified in advance one can find  $a^* \in D, \varepsilon_2$  and  $\delta_1$  for which  $v'(t, x, a^*) \leq 0$  holds together with (2.4) for all  $t \geq t_0$  and  $x \in \bar{\Pi}_\varepsilon$ , then the motion  $x = 0$  possesses stiffness with respect to  $x_\alpha$ .

Proof. Assume that the theorem's hypotheses are fulfilled: the parameter  $a^*$  has been defined and the domains  $\bar{\Pi}_\varepsilon$  and  $\bar{\Pi}_\delta$  have been constructed for arbitrary  $\varepsilon_1, \delta_2$  and  $t_0$ . Then the solution  $x(t, t_0, x_0, a^*) \in \bar{\Pi}_\varepsilon$  if only  $x_0 \in \bar{\Pi}_\delta$ . As a matter of fact, arguing otherwise, we assume that when  $a = a^*$  a solution  $x(t)$  exists reaching the boundary of  $\bar{\Pi}_\varepsilon$  at an instant  $t_1 > t_0$  notwithstanding that the condition  $x_0 \in \bar{\Pi}_\delta$  obtains at  $t = t_0$ . Since solution  $x(t) \in \bar{\Pi}_\varepsilon$  when  $t \in [t_0, t_1]$ , function  $v$  does not grow along it; consequently,  $v(t_1, x(t_1), b^*) \leq v(t_0, x_0, b^*)$ , which contradicts (2.4). The motion  $x = 0$  possesses the stiffness property.

Corollary 1. If a function  $v(t, x, b)$  possessing property (A) with respect to  $x_\alpha$  exists and  $v'(t, x, a) \leq 0$  for all  $t \geq t^*, x \in \Gamma$  and  $a \in D$ , then motion  $x = 0$  possesses stiffness with respect to  $x_\alpha$ .

Theorem 2. If for system (1.4):

- a) there exists a function  $v(t, x, b)$  possessing property (A) with respect to  $x_\alpha$  uniformly in  $t \in [t^*, \infty)$ ;
- b) function  $v$  and its derivative  $v'$  satisfy the condition: for every  $\varepsilon_1$  and  $\delta_2$  specified in advance one can find  $a^* \in D, \varepsilon_2$  and  $\delta_1$  for which  $v'(t, x, a^*) \leq 0$  holds together with (2.6) for all  $t \geq t^*$  and  $x \in \bar{\Pi}_\varepsilon \setminus \Pi_\delta$ , then the motion  $x = 0$  possesses stiffness with respect to  $x_\alpha$  uniformly in  $t_0 \in [t^*, \infty)$ .

Proof. Arguing to the contrary, we assume that when  $a = a^*$  a solution  $x(t)$  ( $x(t_0) = x_0, x_0 \in \bar{\Pi}_\delta, t_0 \geq t^*$ ) exists reaching the boundary of  $\bar{\Pi}_\varepsilon$  at an instant  $t_1 > t_0$ . Let  $t'$  ( $t_0 \leq t' < t_1$ ) be an instant for which  $x(t') \in \bar{\Pi}_\delta \setminus \Pi_\delta$  and let  $x(t) \in \bar{\Pi}_\varepsilon \setminus \Pi_\delta$  if  $t \in [t', t_1]$ . Function  $v$  does not grow along solution  $x(t)$  on the time interval indicated and, therefore,  $v(t_1, x(t_1), b^*) \leq v(t', x(t'), b^*)$ . The latter contradicts condition (2.6). Consequently,  $x(t) \in \bar{\Pi}_\varepsilon$  if  $x_0 \in \bar{\Pi}_\delta$  and  $t_0 \geq t^*$ .

Theorem 3. If for system (1.4):

- a) a function  $v(t, x, b)$  exists possessing property (A) with respect to  $x_\alpha$  and admitting of an infinitesimal upper bound at  $x = 0$ ;
- b) function  $v$  and its derivative  $v'$  satisfy the condition: for every  $\varepsilon_1, \delta_2$  and  $t_0 \geq t^*$  specified in advance we can find  $a^* \in D, \varepsilon_2$  and  $\delta_1$  for which (2.4) is fulfilled and  $v(t, x, b^*)$  is positive definite in domain  $\bar{\Pi}_\varepsilon$  when  $t \geq t_0$  while  $v'(t, x, a^*)$  is negative definite, then the motion  $x = 0$  possesses strong stiffness with respect to  $x_\alpha$ .

Proof. Function  $v$  and its derivative  $v'$  satisfy the hypotheses of Theorem 1 and so the motion  $x = 0$  possesses stiffness with respect to  $x_\alpha$ . Consequently, every solution  $x(t, t_0, x_0, a^*) \in \bar{\Pi}_\varepsilon$  when  $t \geq t_0$  if  $x_0 \in \bar{\Pi}_\delta$ . It remains to show

that  $x(t, t_0, x_0, a^*) \rightarrow 0$  as  $t \rightarrow \infty$ . The latter can be established by using, say, the proof scheme of Theorem II in [7].

**Corollary 2.** If a function  $v(t, x, b)$  exists that possesses property (A) with respect to  $x_\alpha$ , is positive definite and admits of an infinitesimal upper bound at  $x = 0$ , while  $v^*(t, x, a)$  is negative definite for all  $t \geq t^*$ ,  $x \in \Gamma$  and  $a \in D$ , then the motion  $x = 0$  possesses strong stiffness with respect to  $x_\alpha$ .

**Note 2.** Theorem 1 can be extended to practical stability if functions possessing property (B) are used.

3°. We consider system (1.4) under constantly acting perturbations

$$dx/dt = X(t, x, a) + \mu R(t, x, a) \quad (\mu = \text{const} > 0) \tag{2.7}$$

We remark that the need for investigating similar systems with a small parameter  $\mu$  arises, for instance, in the theory of oscillations and in other problems. Besides the usual requirements on the functions  $R_s(t, x, a)$  ( $s = 1, \dots, n$ ), we shall assume their uniform boundedness in each domain  $\bar{\Pi}_\varepsilon \subset \Gamma_0$  from (1.5) when  $t \geq t^*$ .

Motion  $x = 0$  possesses stiffness with respect to  $x_\alpha$  constantly acting perturbations if for any  $\varepsilon_1, \delta_2$  and  $t_0$  we can find  $a^*$  and  $\mu^*$  depending on them and  $\varepsilon_2$  and  $\delta_1$ , defining domain (1.6) for which the solution of (2.7)  $x(t, t_0, x_0, a^*, \mu^*) \in \Pi_\varepsilon (t \geq t_0)$ , if only  $x_0 \in \bar{\Pi}_\delta$ , for any function  $R_s$ .

Let us assume that a function  $v(t, x, b)$  satisfying the hypotheses of Theorem 2 has been constructed for system (1.4), with the following additions:  $v^*(t, x, a^*) < -l$  ( $l = \text{const} > 0$ ),  $x \in \bar{\Pi}_\varepsilon \setminus \Pi_\delta$ ,  $t \geq t^*$ , and the derivatives  $\partial v(t, x, b^*) / \partial x_s$  ( $s = 1, \dots, n$ ) are uniformly bounded in domain  $\bar{\Pi}_\varepsilon$  when  $t \geq t^*$ . Then the motion  $x = 0$  possesses stiffness under constantly acting perturbations.

Indeed, the motion  $x = 0$  of system (1.4) possesses stiffness with respect to  $x_\alpha$  uniformly in  $t_0 \in [t^*, \infty)$ . We assume that parameter  $a^*$  has been fixed and the domains  $\bar{\Pi}_\varepsilon$  and  $\bar{\Pi}_\delta$  for which (2.6) holds have been constructed, and also that  $v^*(t, x, a^*) < -l$  ( $x \in \bar{\Pi}_\varepsilon \setminus \Pi_\delta$ ,  $t \geq t^*$ ). In addition,  $|\partial v(t, x, b^*) / \partial x_s| < N$  and  $|R_s(t, x, a^*)| < M$ . We set up the expression for the derivative of function  $v(t, x, b^*)$  by virtue of system (2.7). We obtain

$$v^*(t, x, a^*, \mu)_{(2.7)} = v^*(t, x, a^*) + \mu \sum_{s=1}^n \frac{\partial v(t, x, b^*)}{\partial x_s} R_s(t, x, a^*)$$

whence it follows that if  $x \in \bar{\Pi}_\varepsilon \setminus \Pi_\delta$ ,  $t \geq t^*$ , then

$$v^*(t, x, a^*, \mu)_{(2.7)} < -l + \mu nNM$$

and  $v^*(t, x, a^*, \mu^*)_{(2.7)} < 0$  for  $\mu^* < l / nNM$ . The latter signifies that the solution of system (2.7) with initial conditions  $x_0 \in \bar{\Pi}_\delta$  and  $t_0 \geq t^*$  do not leave the domain  $\Pi_\varepsilon$  for  $t > t_0$ .

As we can see the theorems presented are in a known sense the analogs of motion stability theorems. Indications of nonstiff motion, which we do not discuss here, can be established in similar fashion.

**Example 2. Stiffness of a vertically rectified Lagrange gyroscope.** For the axis  $z_1$  directed vertically upward the unperturbed motion corresponds to the values (1.7) of the variables, while the perturbed motion, to (1.8). Since variable  $\eta_3$  can be expressed in terms of  $\eta_1$  and  $\eta_2$  (see (1.9)), the equations of perturbed motion, depending on parameter  $r$ , can be written for the variables  $\xi_i$  and  $\eta_i$  ( $i = 1, 2$ ). Using the Lagrange integrals for the equations of motion, we can write the integrals of perturbed motion

$$v_1 = A(\xi_1^2 + \xi_2^2) - 2mgz\eta_3, \quad v_2 = A(\xi_1\eta_1 + \xi_2\eta_2) - Cr\eta_3 \quad (2.8)$$

$$(\eta_3 = 1 - (1 - \eta_1^2 - \eta_2^2)^{1/2})$$

here  $z$  is the coordinate of the center of gravity.

Let us consider the bundle of integrals (2.8)

$$v(x, r) = v_1 - \lambda v_2 = A(\xi_1^2 + \xi_2^2) - A\lambda(\xi_1\eta_1 + \xi_2\eta_2) + (\lambda Cr - 2mgz)\eta_3 \quad (2.9)$$

where  $\lambda$  is a constant not determined as yet. Let us show that the function  $v$  in (2.9) possesses property (A) with respect to  $\eta_1$  and  $\eta_2$ , satisfying the hypotheses of the Lemma presented above. For this purpose we make the required constructions, using domains (1.10).

Let  $\varepsilon_1 < 1$  and  $\delta_2$  be given. Since  $v^*(\xi) = A(\xi_1^2 + \xi_2^2)$ ,

$$M = \sup [v(x, r): \eta_1 = \eta_2 = 0, |\xi_1| \leq \delta_2, |\xi_2| \leq \delta_2] = 2A\delta_2^2$$

On the set  $\eta_1^2 + \eta_2^2 = \varepsilon_1^2$  the function  $v$  has the minimum

$$\min v = (\lambda Cr - 2mgz)(1 - \sqrt{1 - \varepsilon_1^2}) - 1/4 A\lambda^2 \varepsilon_1^2 \quad (2.10)$$

depending on the bundle's parameter  $\lambda$ . Having chosen the magnitude of this parameter from the condition that expression (2.10) be maximum, we obtain

$$\max \min v = C^2 r^2 (1 - \sqrt{1 - \varepsilon_1^2})^2 / A\varepsilon_1^2 - 2mgz(1 - \sqrt{1 - \varepsilon_1^2}) \quad (2.11)$$

We require that  $\max \min v > M$ . In accord with (2.11) we obtain the following condition for choosing the magnitude of parameter  $r$ :

$$C^2 r^2 > 2A\varepsilon_1^2 [A\delta_2^2 + mgz(1 - \sqrt{1 - \varepsilon_1^2})] / (1 - \sqrt{1 - \varepsilon_1^2})^2 \quad (2.12)$$

Function (2.9) satisfies the last condition of the Lemma. Indeed,  $v(x, r) \rightarrow +\infty$  as  $\xi_1^2 + \xi_2^2 \rightarrow \infty$  uniformly in the domain  $\eta_1^2 + \eta_2^2 \leq \varepsilon_1^2$ .

Having now set  $r = \omega + \xi_3$  (2.12), where  $|\xi_3| \leq \delta_2$ , we obtain

$$|\omega| > \delta_2 + \sqrt{2A\varepsilon_1 [A\delta_2^2 + mgz(1 - \sqrt{1 - \varepsilon_1^2})]^{1/2}} / C(1 - \sqrt{1 - \varepsilon_1^2}) \quad (2.13)$$

We note that the numbers  $\delta_1$  and  $\varepsilon_2$  defining for the variables  $\xi_i$  and  $\eta_i$  ( $i = 1, 2$ ) the domains (1.10) depend on  $\xi_3$ . Keeping in mind that  $|\xi_3| \leq \delta_2$ , we can take

$\delta_1 = \inf \delta_1(\xi_3)$  and  $\varepsilon_2 = \sup \varepsilon_2(\xi_3)$ . Thus the function  $v$  of (2.9), being an integral of the equations of perturbed motion, possesses property (A) with respect to  $\eta_1$  and  $\eta_2$ . Then by Corollary 1 we can deduce that motion (1.7) possesses stiffness with respect to these variables.

In concluding the analysis of the example we note that (2.13) becomes (1.13) when  $z = 0$ . It can also be shown that (2.12) is fulfilled for sufficiently small  $\varepsilon_1$  and  $\delta_2$  if the stability condition  $C^2\omega^2 > 4Amgz$  [8] holds. Indeed, assuming that  $\delta_2/\varepsilon_1 \rightarrow 0$  as  $\varepsilon_1 \rightarrow 0$ , we get that the limit of the right-hand side of (2.13) equals  $4Amgz$ .

**Example 3. Property of stiffness-equilibrium of a conservative system.** The equilibrium position of a system subject to holonomic and stationary constraints is determined by the generalized coordinates  $q_i$  ( $i = 1, \dots, n$ ). We consider the case when the system's potential energy  $\Pi = \Pi(q_1, \dots, q_n, a_1, \dots, a_r)$  depends upon parameters and we assume that  $q_i = 0$  is an isolated equilibrium position for each  $a \in D$ . We assume that  $\Pi(0, a) \equiv 0$ . We denote

$$\varphi(\varepsilon, a) = \inf \{ \Pi(q, a) : q \in K_\varepsilon \setminus \bar{K}_\varepsilon \}, \quad \bar{K}_\varepsilon = \{q : |q_i| \leq \varepsilon\}$$

where  $\varepsilon$  is sufficiently small.

The equilibrium position  $q = 0$  possesses stiffness with respect to the coordinates if for any positive  $\varepsilon$  and  $N$  (the first arbitrarily small, the second, arbitrarily large) a parameter  $a^* \in D$  depending on them exists for which  $\varphi(\varepsilon, a^*) > N$ . Indeed, under the assumptions made the system's total energy

$$v(q, q', a) = \frac{1}{2} \sum_{i,j=1}^n a_{ij}(q) q_i' q_j' + \Pi(q, a)$$

possesses property (A) with respect to  $q$ , having satisfied the Lemma's conditions. Since by virtue of the equations of motion  $v' \equiv 0$ , integral  $v$  satisfies the hypotheses of Corollary 1.

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